

# Charge Duality in Off-Shell Electromagnetism<sup>\*</sup>

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<sup>\*</sup> Stueckelberg-Schrodinger Relativistic Quantum Theory and  
its Associated 5D Electromagnetic Theory

# Off-Shell Electromagnetic Fields

Field equations

$$\partial_\beta f^{\alpha\beta} = ej^\alpha \quad \epsilon^{\alpha\beta\gamma\delta\epsilon} \partial_\alpha f_{\beta\gamma} = 0$$

$$f_{\alpha\beta}(x, \tau) = \partial_\alpha a_\beta(x, \tau) - \partial_\beta a_\alpha(x, \tau)$$

Flat metric

$$g_{\alpha\beta} = \text{diag}(-1, 1, 1, 1, \sigma), \quad \mu, \nu, \rho = 0, \dots, 3 \quad x^5 = \tau$$

$$\alpha, \beta, \gamma = 0, \dots, 3, 5$$

Connection with  
Maxwell theory

$$\rho \xrightarrow{\tau \rightarrow \pm\infty} 0 \Rightarrow \partial_\mu J^\mu(x) = \partial_\mu \int_{-\infty}^{\infty} d\tau j^\mu(x, \tau) = 0$$

$$f^{5\mu} \xrightarrow{\tau \rightarrow \pm\infty} 0 \Rightarrow \left. \begin{array}{l} \partial_\beta f^{\alpha\beta} = ej^\alpha \\ \epsilon^{\alpha\beta\gamma\delta\epsilon} \partial_\alpha f_{\beta\gamma} = 0 \end{array} \right\} \rightarrow \left\{ \begin{array}{l} \partial_\nu F^{\mu\nu} = eJ^\mu \\ \epsilon^{\mu\nu\rho\lambda} \partial_\mu F_{\nu\rho} = 0 \end{array} \right.$$

$$F^{\mu\nu}(x) = \int_{-\infty}^{\infty} d\tau f^{\mu\nu}(x, \tau)$$

$$A^\mu(x) = \int_{-\infty}^{\infty} d\tau a^\mu(x, \tau)$$

# Space-Time (Clifford) Algebra

Vector  
Product

$$ab = \frac{1}{2}(ab + ba) + \frac{1}{2}(ab - ba) = a \cdot b + a \wedge b$$

Multivector

$$\begin{aligned} A &= A_0 + A_1 + A_2 + A_3 + \cdots + A_D \\ &= A_0 + A_1^i e_i + A_2^{ij} e_i \wedge e_j + \cdots + A_D^{i_0 i_2 \cdots i_{D-1}} e_{i_0} \wedge e_{i_1} \wedge \cdots \wedge e_{i_{D-1}} \end{aligned}$$

Basis:  $2^n$   
dimensions

$$\{1, e_i, e_i \wedge e_j, e_i \wedge e_j \wedge e_k, \dots, e_0 \wedge e_1 \wedge \cdots \wedge e_{D-1}\}$$

Multivector  
Products

$$\begin{aligned} aA_r &= a(a_1 \wedge a_2 \wedge \cdots \wedge a_r) = a \cdot A_r + a \wedge A_r \\ a \cdot A_r &= \frac{1}{2}[aA_r - (-1)^r A_r a] \\ &= \sum_{i=1}^r (-1)^{i+1} (a \cdot a_i) a_1 \wedge \cdots \wedge a_{i-1} \wedge a_{i+1} \wedge \cdots \wedge a_r \\ a \wedge A_r &= \frac{1}{2}[aA_r + (-1)^r A_r a] = a \wedge a_1 \wedge a_2 \wedge \cdots \wedge a_r \end{aligned}$$

## Relationships Among Products

$$\mathbf{i} = \mathbf{e}_0 \wedge \mathbf{e}_1 \wedge \cdots \wedge \mathbf{e}_{D-1} \qquad \mathbf{i}^2 = (-1)^{\frac{D(D-1)}{2}} g_{00} \cdots g_{D-1,D-1}$$

Pseudoscalar

$$\mathbf{i} \left[ \mathbf{e}_{i_1} \wedge \cdots \wedge \mathbf{e}_{i_n} \right] = g_{i_1 i_1} \cdots g_{i_n i_n} \frac{1}{n!} \sum_{i_1 \cdots i_n} \epsilon_{i_1 \cdots i_n}^{i_{n+1} \cdots i_D} \left[ \mathbf{e}_{i_{n+1}} \wedge \cdots \wedge \mathbf{e}_{i_D} \right]$$

Symmetry

$$\mathbf{e}_k \mathbf{i} = \mathbf{e}_k \mathbf{e}_0 \mathbf{e}_1 \cdots \mathbf{e}_k \cdots \mathbf{e}_{D-1} = (-1)^{k-1} \mathbf{e}_0 \mathbf{e}_1 \cdots \mathbf{e}_k \mathbf{e}_k \cdots \mathbf{e}_{D-1} = (-1)^{D-1} \mathbf{i} \mathbf{e}_k$$

$$\mathbf{i} (a \wedge B_r) = \frac{1}{2} [\mathbf{i} a B_r + (-1)^r \mathbf{i} B_r a]$$

Useful Identity

$$\begin{aligned} a \cdot (\mathbf{i} B_r) &= \frac{1}{2} [a \mathbf{i} B_r - (-1)^{D-r} \mathbf{i} B_r a] \\ &= \frac{1}{2} [(-1)^{D-1} \mathbf{i} a B_r - (-1)^{D-r} \mathbf{i} B_r a] \\ &= (-1)^{D-1} \frac{1}{2} [\mathbf{i} a B_r + (-1)^r \mathbf{i} B_r a] \\ &= (-1)^{D-1} \mathbf{i} (a \wedge B_r) \end{aligned}$$

# Formulation of Field Equations

Definitions

$$\mathbf{d} = \partial^\alpha \mathbf{e}_\alpha \quad \mathbf{f} = 1/2 f^{\alpha\beta} \mathbf{e}_\alpha \wedge \mathbf{e}_\beta$$

$$\mathbf{j}_e = j^\alpha \mathbf{e}_\alpha \quad \mathbf{e}_\alpha \cdot \mathbf{e}_\beta = g_{\alpha\beta}$$

Field Equations

$$d\mathbf{f} = -e\mathbf{j}_e \Rightarrow \mathbf{d} \cdot \mathbf{f} + \mathbf{d} \wedge \mathbf{f} = -e\mathbf{j}_e$$

$$\begin{aligned} \mathbf{d} \cdot \mathbf{f} = -e\mathbf{j}_e \quad \rightarrow \quad \frac{1}{2} (\partial^\gamma \mathbf{e}_\gamma) \cdot (f^{\alpha\beta} \mathbf{e}_\alpha \wedge \mathbf{e}_\beta) &= \frac{1}{2} \partial^\gamma f^{\alpha\beta} \mathbf{e}_\gamma \cdot (\mathbf{e}_\alpha \wedge \mathbf{e}_\beta) \\ &= \frac{1}{2} \partial^\gamma f^{\alpha\beta} (g_{\gamma\alpha} \mathbf{e}_\beta - g_{\gamma\beta} \mathbf{e}_\alpha) \\ &= \partial_\alpha f^{\alpha\beta} \mathbf{e}_\beta = j^\beta \mathbf{e}_\beta \end{aligned}$$

$$\begin{aligned} \mathbf{d} \wedge \mathbf{f} = 0 \quad \rightarrow \quad i\mathbf{d} \wedge \mathbf{f} = \mathbf{d} \cdot i\mathbf{f} &= (\partial^\alpha \mathbf{e}_\alpha) \cdot \left[ 1/(D-2)! \epsilon^{\alpha_1 \alpha_2 \dots \alpha_{D-2} \beta \gamma} f_{\beta\gamma} \mathbf{e}_{\alpha_1} \mathbf{e}_{\alpha_2} \dots \mathbf{e}_{\alpha_{D-2}} \right] \\ &= \left[ 1/(D-2)! \epsilon^{\alpha_1 \alpha_2 \dots \alpha_{D-3} \alpha \beta \gamma} \partial_\alpha f_{\beta\gamma} \mathbf{e}_{\alpha_1} \mathbf{e}_{\alpha_2} \dots \mathbf{e}_{\alpha_{D-3}} \right] = 0 \end{aligned}$$

D = 4

$$\partial_\nu F^{\mu\nu} = eJ^\mu \quad \partial_\mu \epsilon^{\mu\nu\rho\lambda} F_{\nu\rho} = 0$$

D = 5

$$\partial_\beta f^{\alpha\beta} = e j^\alpha \quad \partial_\alpha \epsilon^{\alpha\beta\gamma\delta\epsilon} f_{\beta\gamma} = 0$$

# Duality and the Dirac Monopole

Suppose a magnetic current

$$\nabla \cdot \mathbf{E} = eJ_e^0$$

$$\nabla \times \mathbf{H} - \partial_0 \mathbf{E} = e\mathbf{J}_e$$

$$\nabla \cdot \mathbf{H} = gJ_m^0$$

$$\nabla \times \mathbf{E} + \partial_0 \mathbf{H} = g\mathbf{J}_m$$

Perform duality rotation

$$\mathbf{E} = \mathbf{E}' \cos \theta + \mathbf{H}' \sin \theta$$

$$\mathbf{H} = -\mathbf{E}' \sin \theta + \mathbf{H}' \cos \theta$$

$$e\mathbf{J}_e = e\mathbf{J}_e' \cos \theta + g\mathbf{J}_m' \sin \theta$$

$$eJ_e^0 = eJ_e'^0 \cos \theta + gJ_m'^0 \sin \theta$$

$$g\mathbf{J}_m = -e\mathbf{J}_e' \sin \theta + g\mathbf{J}_m' \cos \theta$$

$$eJ_e^0 = -eJ_e'^0 \sin \theta + gJ_m'^0 \cos \theta$$

Choose angle so that

$$\nabla \cdot \mathbf{E}' = eJ_e^0$$

$$\nabla \times \mathbf{H}' - \partial_0 \mathbf{E}' = e\mathbf{J}_e'$$

$$\nabla \cdot \mathbf{H}' = 0$$

$$\nabla \times \mathbf{E}' + \partial_0 \mathbf{H}' = 0$$

Quantization condition

$$\frac{ge}{\hbar c} = \frac{n}{2}$$

Magnetic monopole  $\Rightarrow$  Charge quantization

Equivalent duality rotation

$$\epsilon^{\mu\nu\lambda\rho} F_{\lambda\rho}(\mathbf{E}, \mathbf{H}) = F^{\mu\nu}(-\mathbf{H}, \mathbf{E})$$

$$F^{\mu\nu} \rightarrow F^{\mu\nu} \cos \theta + \epsilon^{\mu\nu\lambda\rho} F_{\lambda\rho} \sin \theta$$

$$\epsilon^{\mu\nu\lambda\rho} F^{\mu\nu} \rightarrow -F^{\mu\nu} \sin \theta + \epsilon^{\mu\nu\lambda\rho} F_{\lambda\rho} \cos \theta$$

$$eJ_e^\mu \rightarrow eJ_e^\mu \cos \theta + gJ_m^\mu \sin \theta$$

$$gJ_m^\mu = -eJ_e^\mu \sin \theta + gJ_m^\mu \cos \theta$$

# Duality Transformation

General field equations  $df = j = -ej_1 - gj_3$

Electric current equation  $d \cdot f = -ej_1$

Magnetic current equation  $d \wedge f = -gj_3$

$d \wedge ih = id \cdot h = -giq_{D-3} \rightarrow d \cdot h = -gq_{D-3}$

1-parameter transformation

$$U = e^{\theta G} = 1 + \theta G + o(\theta^2)$$

$$Udf = Uj$$

$$(1 + \theta G)(d \cdot f + d \wedge f) = (1 + \theta G)(-ej_1 - gj_3)$$

$$d \cdot f + d \wedge f + \theta G(d \cdot f) + \theta G(d \wedge f) = (-ej_1 - gj_3 - e\theta Gj_1 - g\theta Gj_3)$$

Require separation by rank

Vector equation  $d \cdot f + \theta G(d \wedge f) = (-ej_1 - g\theta Gj_3)$

Trivector equation  $d \wedge f + \theta G(d \cdot f) = (-gj_3 - e\theta Gj_1)$

# Duality Symmetry

Require  
form  
invariance

$$\left. \begin{array}{l} G(\mathbf{d} \wedge \mathbf{f}) = \mathbf{d} \cdot (G\mathbf{f}) \\ G(\mathbf{d} \cdot \mathbf{f}) = \mathbf{d} \wedge (G\mathbf{f}) \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} \mathbf{d} \cdot (\mathbf{f} + \theta G\mathbf{f}) = -e\mathbf{j}_1 - g\theta G\mathbf{j}_3 \\ \mathbf{d} \wedge (\mathbf{f} + \theta G\mathbf{f}) = -g\mathbf{j}_3 - e\theta G\mathbf{j}_1 \end{array} \right.$$

Requirements  
on  $G$

$\mathbf{j}_1$  and  $G\mathbf{j}_3$  must be vectors ( $G : 3 \rightarrow 1$ )

$\mathbf{j}_3$  and  $G\mathbf{j}_1$  must be trivectors ( $G : 1 \rightarrow 3$ )

$G\mathbf{f}$  must be a bivector ( $G : 2 \rightarrow 2$ )

$$G(\mathbf{d} \cdot \mathbf{f}) = \mathbf{d} \wedge (G\mathbf{f}) \Rightarrow G\mathbf{d} = -\mathbf{d}G$$

Requirements are only satisfied for  $G = -i$  in  $D = 4$

Choose  
 $\theta$  so  
that

$$\left. \begin{array}{l} \mathbf{d} \cdot (\mathbf{f} - \theta i\mathbf{f}) = \mathbf{d} \cdot \mathbf{f}' \\ -e\mathbf{j}_1 - g\theta i\mathbf{j}_3 = -e\mathbf{j}_1' \\ \mathbf{d} \wedge (\mathbf{f} - \theta i\mathbf{f}) = \mathbf{d} \wedge \mathbf{f}' \\ -g\mathbf{j}_3 - e\theta i\mathbf{j}_1 = 0 \end{array} \right\} \rightarrow \left. \begin{array}{l} \mathbf{d} \cdot \mathbf{f}' = -e\mathbf{j}_1' \\ \mathbf{d} \wedge \mathbf{f}' = 0 \end{array} \right\} \left| \begin{array}{l} U = e^{-i\theta} \\ U\mathbf{d}\mathbf{f} = \mathbf{d}U\mathbf{f} = \mathbf{d}\mathbf{f}' \\ U\mathbf{j} = U(-e\mathbf{j}_1 - g\mathbf{j}_3) = -e\mathbf{j}_1' \end{array} \right.$$

# Generalized Electromagnetic Fields in D=5

Multivector  
fields and  
currents

$$dF = J$$

$$F = F_0 + F_1 + F_2 + F_3 + F_4 + F_5$$

$$J = J_0 + J_1 + J_2 + J_3 + J_4 + J_5$$

Separating  
Terms

$$d \cdot F_1 = J_0 \quad d \wedge F_2 + d \cdot F_4 = J_3$$

$$dF_0 + d \cdot F_2 = J_1 \quad d \wedge F_3 + d \cdot F_5 = J_4$$

$$d \wedge F_1 + d \cdot F_3 = J_2 \quad d \wedge F_4 = J_5$$

Duality  
Transformation

$$U = 1 + \theta i + o(\theta^2)$$

$$d \cdot F_1 + \theta d \cdot iF_4 = J_0 + \theta iJ_5$$

$$dF_0 + d \cdot F_2 + \theta d \cdot iF_3 + \theta id \cdot F_5 = J_1 + \theta iJ_4$$

$$d \wedge F_1 + d \cdot F_3 + \theta d \cdot iF_2 + \theta d \wedge iF_4 = J_2 + \theta iJ_3$$

$$d \wedge F_2 + d \cdot F_4 + \theta d \cdot iF_1 + \theta d \wedge iF_3 = J_3 + \theta iJ_2$$

$$d \wedge F_3 + d \cdot F_5 + \theta d \wedge iF_2 = J_4 + \theta iJ_1$$

$$d \wedge F_4 + \theta idF_0 + \theta d \wedge iF_1 = J_5 + \theta iJ_0$$

## Identify Associated Terms

$$d \cdot F_1 + \theta d \cdot iF_4 = J_0 + \theta iJ_5 \rightarrow d \cdot (F_1 + \theta iF_4) = \boxed{d \cdot \tilde{F}_1 = \tilde{J}_0}$$

$$d \cdot F_2 + \theta d \cdot iF_3 + \theta id \cdot F_5 = J_1 + \theta iJ_4 \rightarrow d \cdot (F_2 + \theta iF_3) = \boxed{d \cdot \tilde{F}_2 = \tilde{J}_1}$$

$$d \wedge F_1 + d \cdot F_3 + \theta d \cdot iF_2 + \theta d \wedge iF_4 = J_2 + \theta iJ_3$$

$$d \wedge (F_1 + \theta iF_4) + d \cdot (F_3 + \theta iF_2) = \boxed{d \wedge \tilde{F}_1 + d \cdot \tilde{F}_3 = \tilde{J}_2}$$

$$d \wedge F_2 + d \cdot F_4 + \theta d \cdot iF_1 + \theta d \wedge iF_3 = J_3 + \theta iJ_2$$

$$d \wedge (F_2 + \theta iF_3) + d \cdot (F_4 + \theta iF_1) = \boxed{d \wedge \tilde{F}_2 + d \cdot \tilde{F}_4 = \tilde{J}_3}$$

$$d \wedge F_3 + \theta d \wedge iF_2 = J_4 + \theta iJ_1 \rightarrow d \wedge (F_3 + \theta iF_2) = \boxed{d \wedge \tilde{F}_3 = \tilde{J}_4}$$

$$d \wedge F_4 + \theta d \wedge iF_1 = J_5 + \theta iJ_0 \rightarrow d \wedge (F_4 + \theta iF_1) = \boxed{d \wedge \tilde{F}_4 = \tilde{J}_5}$$

## Minimum Set of Fields

Generalized  
field  
equations

$$d \cdot \tilde{F}_1 = \tilde{J}_0$$

$$d \wedge \tilde{F}_2 + d \cdot \tilde{F}_4 = \tilde{J}_3$$

$$d \cdot \tilde{F}_2 = \tilde{J}_1$$

$$d \wedge \tilde{F}_3 = \tilde{J}_4$$

$$d \wedge \tilde{F}_1 + d \cdot \tilde{F}_3 = \tilde{J}_2$$

$$d \wedge \tilde{F}_4 = \tilde{J}_5$$

Eliminate scalar and  
pseudoscalar current  
sectors

$$d \cdot \tilde{F}_2 = \tilde{J}_1$$

$$d \wedge \tilde{F}_2 = \tilde{J}_3 = 0$$

$$d \cdot \tilde{F}_3 = \tilde{J}_2$$

$$d \wedge \tilde{F}_3 = \tilde{J}_4$$

Set magnetic current  
to zero

Standard field equations  
without magnetic current

$$d \cdot \tilde{F}_2 = \tilde{J}_1$$

$$d \wedge \tilde{F}_2 = 0$$

Equivalent field equations  
with magnetic current

$$d \cdot \tilde{F}_3 = d \cdot iH_2 = \tilde{J}_2 \rightarrow d \wedge H_2 = i\tilde{J}_2 = Q_3$$

$$d \cdot H_2 = i\tilde{J}_4 = Q_1$$

## Full Set of Fields

Generalized  
field equations

$$d \cdot \tilde{F}_1 = \tilde{J}_0$$

$$d \cdot \tilde{F}_2 = \tilde{J}_1$$

$$d \wedge \tilde{F}_1 + d \cdot \tilde{F}_3 = \tilde{J}_2$$

$$d \wedge \tilde{F}_2 + d \cdot \tilde{F}_4 = \tilde{J}_3$$

$$d \wedge \tilde{F}_3 = \tilde{J}_4$$

$$d \wedge \tilde{F}_4 = \tilde{J}_5$$

2-field sector without  
magnetic current

$$d \cdot \tilde{F}_2 = \tilde{J}_1$$

$$d \wedge \tilde{F}_2 = 0$$

4-field

$$d \cdot \tilde{F}_4 = \tilde{J}_3 \rightarrow d \wedge H_1 = Q_2$$

$$d \wedge \tilde{F}_4 = \tilde{J}_5 \rightarrow d \cdot H_1 = Q_0$$

3-field  
and  
1-field

$$d \wedge \tilde{F}_1 + d \cdot \tilde{F}_3 = \tilde{J}_2 \rightarrow d \cdot H_4 + d \wedge H_2 = Q_3$$

$$d \wedge \tilde{F}_3 = \tilde{J}_4 \rightarrow d \cdot H_2 = Q_1 \quad d \cdot \tilde{F}_1 = \tilde{J}_0 \rightarrow d \wedge H_4 = Q_5$$

Equivalent  
2-field sector with  
magnetic current

$$d \cdot H_2 = Q_1 \quad d \cdot H_4 + d \wedge H_2 = Q_3 \quad d \wedge H_4 = Q_5$$

Equivalent  
1-field sector

$$d \cdot H_1 = Q_0 \quad d \wedge H_1 = Q_2$$

# Lorentz Force With Magnetic Charges

$$\mathbf{F} = e[\mathbf{E} + \mathbf{v} \times \mathbf{H}] + g[\mathbf{H} - \mathbf{v} \times \mathbf{E}]$$

$$M \frac{du^i}{dt} = e[E^i + \epsilon^{ijk} v_j \times H_k] + g[H^i - \epsilon^{ijk} v_j \times E_k]$$

Modified  
Lorentz  
force in  
D=4

$$\begin{aligned} M \frac{du^i}{dt} \frac{dt}{d\tau} &= e \left[ E^i \frac{dt}{d\tau} + \epsilon^{ijk} u_j \times H_k \right] + g \left[ H^i \frac{dt}{d\tau} - \epsilon^{ijk} u_j \times E_k \right] \\ &= e [F^{i0} \dot{x}_0 + F^{ij} u_j] + g \epsilon^{0i\mu\nu} [F_{\mu\nu} \dot{x}_0 + F_{0\mu} \dot{x}_\nu] \end{aligned}$$

$$\begin{aligned} M \frac{d^2 x^i}{d\tau^2} &= e F^{i\mu} \dot{x}_\mu + g \epsilon^{i\mu\nu\lambda} F_{\mu\nu} \dot{x}_\lambda \\ &= e F^{i\mu} \dot{x}_\mu + g \tilde{F}^{i\mu} \dot{x}_\mu \xrightarrow{\text{duality rotation}} e (F')^{i\mu} \dot{x}_\mu \end{aligned}$$

no duality  
rotation in  
D=5

$$\begin{aligned} M \frac{d^2 x^\mu}{d\tau^2} &= e f^{\mu\alpha} \dot{x}_\alpha + g \tilde{f}^{\mu\alpha} \dot{x}_\alpha \\ &= e f^{\mu\alpha} \dot{x}_\alpha + g \epsilon^{\mu\alpha\beta\gamma\delta} f_{\alpha\beta\gamma} \dot{x}_\delta \end{aligned}$$